

Unveiling the complete dispersion of the giant Rashba split surface states of ferroelectric α -GeTe(111) by alkali doping

G. Kremer^{1,*}, T. Jaouen¹, B. Salzmann¹, L. Nicolai², M. Rumo¹, C. W. Nicholson¹, B. Hildebrand¹, J. H. Dil^{3,4}, J. Minár², G. Springholz⁵, J. Krempaský⁴ and C. Monney¹

¹*Département de Physique and Fribourg Center for Nanomaterials, Université de Fribourg, CH-1700 Fribourg, Switzerland*

²*New Technologies-Research Center, University of West Bohemia, Plzen, Czech Republic*

³*Institute of Physics, Ecole Polytechnique Fédérale de Lausanne, CH-1015 Lausanne, Switzerland*

⁴*Photon Science Division, Paul Scherrer Institut, CH-5232 Villigen, Switzerland*

⁵*Institut für Halbleiter- und Festkörperphysik, Johannes Kepler Universität, A-4040 Linz, Austria*



(Received 3 April 2020; revised 3 June 2020; accepted 17 June 2020; published 22 July 2020)

α -GeTe(111) is a noncentrosymmetric ferroelectric material for which a strong spin-orbit interaction gives rise to giant Rashba split states in the bulk and at the surface. The detailed dispersions of the surface states inside the bulk band gap remains an open question because they are located in the unoccupied part of the electronic structure, making them inaccessible to static angle-resolved photoemission spectroscopy. We show that this difficulty can be overcome via *in situ* potassium doping of the surface, leading to a rigid shift of 80 meV of the surface states into the occupied states. Thus, we resolve, in great detail, their dispersion and highlight their crossing at the $\bar{\Gamma}$ point, which, in comparison with density functional theory calculations, definitively confirms the Rashba mechanism.

DOI: [10.1103/PhysRevResearch.2.033115](https://doi.org/10.1103/PhysRevResearch.2.033115)

I. INTRODUCTION

The electronic band structure in solids with both inversion and time-reversal symmetries (TRSS) is, at least, doubly degenerate with respect to the spin degree of freedom. In systems where the inversion symmetry is broken, this degeneracy can be lifted by the spin-orbit interaction (SOI). Bychkov and Rashba have theoretically described this phenomenon by considering a two-dimensional (2D) electron gas subject to an external out-of-plane electric field [1]. The resulting energy levels are split by a linear term in momentum \mathbf{k} : $E^\pm(\mathbf{k}) = E_0 + (\hbar^2 \mathbf{k}^2 / 2m^*) \pm \alpha_R |\mathbf{k}|$, where m^* is the effective mass of the electrons and α_R is the Rashba parameter reflecting the magnitude of this so-called Rashba effect. The Rashba effect was first observed by angle-resolved photoemission spectroscopy (ARPES) at the surface of Au(111) where the Shockley state shows a momentum-dependent energy splitting [2]. Subsequently, spin-resolved ARPES confirmed its predicted in-plane spin polarization [3]. Similar experimental investigations on other surfaces and interfaces with variable splitting magnitude have been performed [4–19]. More recently, work has been done on BiTeX ($X = \text{I, Cl, and Br}$) polar semiconductors where the first evidence of giant bulk Rashba split states was observed by ARPES and spin-resolved ARPES [20–26].

In recent decades, an increasing amount of research has been carried out aiming to enhance the control of spin-polarized currents in nanostructured materials. Rashba systems have appeared as an ideal playground to address these concepts and have been recently used as an efficient spin-to-charge converter in prototype spintronics devices [27–29]. The synthesis of new materials with giant and tunable Rashba split electronic states, adjustable position of the chemical potential and, ideally, spin-polarized states crossing the Fermi level (E_F), is highly desirable in the view of future applications in the growing area of spintronics [30].

In this context, α -GeTe(111) is a promising compound [31]. It is a ferroelectric semiconductor which exhibits spin-polarized bulk and surface split electronic states with the largest Rashba parameter currently reported and has been termed a ferroelectric Rashba semiconductor [32,33]. Bulk and surface states are Rashba split due to out-of-plane inversion symmetry breaking and large SOI.

Considerable work has been carried out to experimentally characterize the electronic band structure of α -GeTe(111), in particular, its bulk Rashba split states [34–36]. Spin-resolved ARPES measurements have experimentally confirmed the link between the spin texture and the ferroelectricity, demonstrating that the spin-polarized band structure can be reversibly manipulated either by modifying the surface termination [37] or with an electric field [38]. Since ARPES experiments can only probe the occupied part of the band structure, these studies have mainly focused on the bulk states, and the surface states still lack a complete and direct experimental characterization. In particular, their detailed dispersion at the $\bar{\Gamma}$ point of the Brillouin zone (BZ) remains unknown because the states are located in the unoccupied part of the band structure. Direct measurements of the surface states is

*Corresponding author: geoffroy.kremer@unifr.ch

Published by the American Physical Society under the terms of the [Creative Commons Attribution 4.0 International](https://creativecommons.org/licenses/by/4.0/) license. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI.

an important point to address in order to fully understand the electronic band structure of the system and to test the validity of current theories.

In the present paper, we show that the surface states of α -GeTe(111) can be shifted down to E_F and experimentally addressed with ARPES via *in situ* surface deposition of potassium (K). Their complete dispersion is resolved by taking advantage of the electronic thermal occupation at room temperature (RT). In particular, we resolve their crossing point (CP) at $\bar{\Gamma}$, thus, confirming the TRS conservation and the Rashba picture. We further demonstrate their nonparabolic dispersion, in excellent agreement with density functional theory (DFT) band-structure calculations. We show that the Rashba-type splitting is unaffected by the electron doping, allowing us to extrapolate the energy position of the surface states on the bare α -GeTe(111).

II. METHODS

Two hundred-nanometers-thick α -GeTe(111) films were grown by molecular beam epitaxy (MBE) on InP(111) substrates and then characterized in a different experimental setup coupling a surface preparation chamber with low-temperature scanning tunneling microscopy (LT-STM), x-ray photoemission spectroscopy (XPS), and ARPES techniques. To avoid surface degradation and oxidation during the transfer between the two systems, a protective stack of amorphous Te- and Se-capping layers was deposited *in situ* after MBE growth and, subsequently, removed in ultrahigh vacuum in the STM/XPS/ARPES system in a base pressure of 7×10^{-11} mbar at 500 K. Note that in contrast to $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Te}_3$ [39] for α -GeTe(111), a pure Te cap does not provide satisfactory results.

The low-temperature (4.5 K) STM images were obtained using an Omicron LT-STM in a pressure better than 5×10^{-11} mbar. The low-energy electron-diffraction (LEED) patterns were recorded with a SPECS ErLEED at RT and 80 eV. ARPES measurements were performed using a Scienta DA30 photoelectron analyzer with monochromatized HeI_α radiation ($h\nu = 21.2$ eV, SPECS UVLS with a TMM 304 monochromator) and at RT, if not further specified. Energy and angular resolutions were better than 10 meV and 0.1° , respectively. XPS measurements were carried out with a monochromatized $\text{AlK}\alpha$ source (SPECS μ -FOCUS, resolution better than 300 meV). The K evaporations were performed at RT from a SAES getter source in a pressure better than 5×10^{-10} mbar.

Ab initio calculations are based on DFT as implemented in the fully relativistic spin-polarized relativistic Korringa-Kohn-Rostoker (KKR) theory [40]. The electronic structure of a semi-infinite surface of α -GeTe is described, including all relativistic effects, by the Dirac equation, which is solved using the screened KKR formalism. The local density approximation based potentials were treated within the atomic sphere approximation and for the multipole expansion of the Green's function an angular momentum cutoff $l_{\text{max}} = 3$ was used. The structural of Te-terminated surface has been taken from the structural relaxation presented in the work of Krempaský *et al.* [38]. The electronic structure of semi-infinite α -GeTe is represented by means of Bloch spectral function (BSF), i.e., imaginary part of the KKR Green's function.

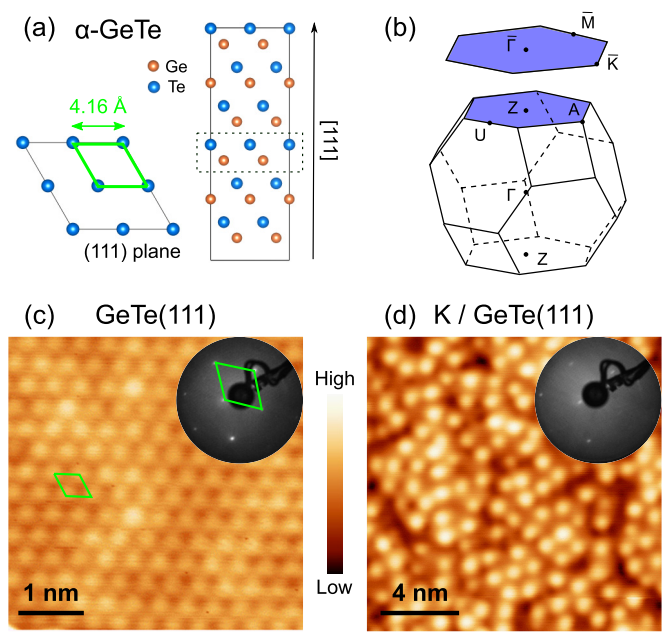


FIG. 1. (a) Crystal structure of α -GeTe along the [111] crystallographic direction (right) and the corresponding projection on the (111) atomic plane (bottom left). The 2D unit-cell (green rhombus) parameter is extracted from x-ray-diffraction measurements [41]. (b) Corresponding three-dimensional BZ and its 2D projection on the (111) plane (blue). (c) STM image of the bare α -GeTe(111) surface ($U = -1.5$ V, $I = 0.15$ nA) and associated LEED pattern in the top right corner. The 2D unit cells in real and reciprocal spaces are indicated in green. (d) STM image of the α -GeTe(111) surface covered by 0.09 monolayer (ML) of K ($U = -1.5$ V, $I = 0.15$ nA) and the associated LEED pattern in the top right corner.

III. RESULTS

The crystal structure of α -GeTe(111) is shown in the right part of Fig. 1(a) (space-group $R3m$). It corresponds to a stacking of “bilayers,” i.e., sequences of Te and Ge planes (see black-dashed rectangle) along the rhombohedral [111] crystallographic direction. The distance between the planes of Te and Ge within the same bilayer is smaller than the one between two bilayers. This large rhombohedral lattice distortion, which is equal to about 10% of the lattice parameter [36], provides the ferroelectric order. The (111) cut of α -GeTe corresponds either to a Te- or a Ge-terminated plane, leading either to an upward or a downward ferroelectric polarization at the surface. In the following, we address the band dispersion of the surface states of a Te-terminated α -GeTe(111) surface, which is predicted to be energetically most favorable [42].

To this end, we have annealed our capped sample at 500 K in order to recover a fresh surface. The resulting surface has first been characterized by STM and LEED as shown in Fig. 1(c). This panel shows an atomically-resolved STM image of the (111) surface of α -GeTe. It reveals hexagonally well-arranged atoms with an interatomic distance estimated as (4.1 ± 0.2) Å, in good agreement with the STM image simulated by DFT calculations of a pristine Te-terminated surface of α -GeTe(111) in the work of Deringer *et al.* [42]. Furthermore, the LEED measurements exhibit a sharp

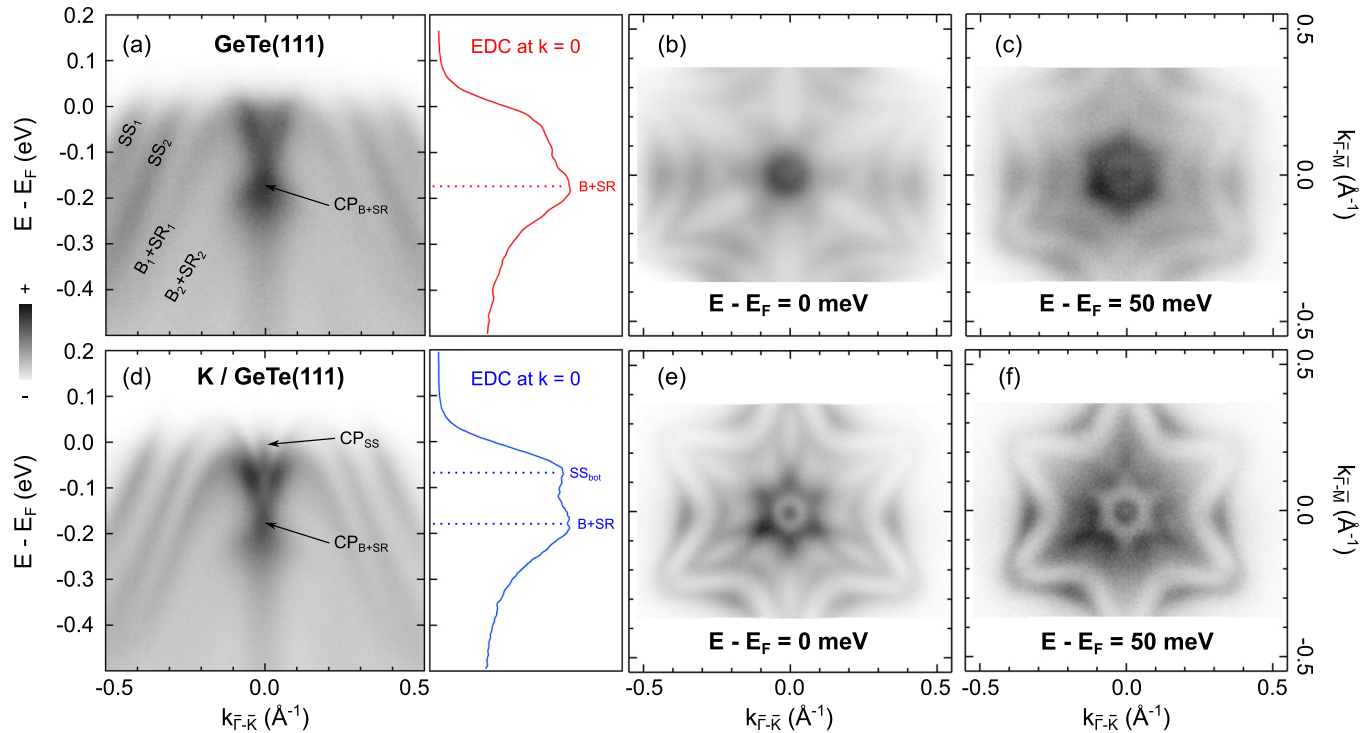


FIG. 2. ARPES spectra along the \bar{K} - $\bar{\Gamma}$ - \bar{K} high-symmetry line of the 2D BZ (left) and the corresponding energy-distribution curve (EDC) taken at $k = 0$ (right) for (a) bare α -GeTe(111) and (d) K/α -GeTe(111). Associated constant energy surfaces taken at $E - E_F = 0$ and $E - E_F = 50$ meV for (b) and (c) bare α -GeTe(111) and (e) and (f) K/α -GeTe(111).

(1×1) hexagonal pattern, demonstrating the long-range surface order and the expected crystal symmetry after the desorption procedure. The chemical composition of our Te-terminated α -GeTe(111) surface was also cross-checked by XPS measurements that show negligible surface contaminants, sharp Ge $2p$ and Te $3d$ core levels, and only residual traces of the Se-capping layer remaining from the decapping procedure (see Fig. 4 in the Appendix).

The electronic band structure of α -GeTe(111) measured by ARPES is shown in Figs. 2(a)–2(c). ARPES intensity along the $\bar{\Gamma}$ - \bar{K} direction [see the 2D BZ in Fig. 1(b)] exhibits all the spectroscopic signatures of a clean Te-terminated surface. Due to spectral broadening, the bulk states B_1 and B_2 are smeared out with their surface resonances SR_1 and SR_2 [36,43]. We also resolve the surface states SS_1 and SS_2 , forming characteristic “snowflake” constant energy surfaces at $E - E_F = 0$ meV [Fig. 2(b)] and $E - E_F = 50$ meV [Fig. 2(c)], in very good agreement with measurements and calculations for Te-terminated α -GeTe(111) [36,37,43]. Based on our STM, LEED, XPS, and ARPES results, we conclude that our α -GeTe(111) surface is very well ordered and Te terminated, both on the atomic and on the macroscopic scales.

Subsequently, we have deposited K atoms on the α -GeTe(111) surface. Figure 1(d) displays the corresponding STM image where the K adatoms are clearly distinguishable as large and bright protrusions. The coverage is estimated as (0.09 ± 0.01) ML (1 ML corresponds to one K atom per α -GeTe(111) surface unit cell). A (1×1) hexagonal pattern is still visible in LEED but with a more-pronounced diffuse

background. This background is due to the randomly organized K atoms at the surface without any coherent surface reconstruction as visible in the STM image. The (1×1) pattern in LEED is due to the α -GeTe(111) sublayer, which is not distinguishable in the STM topography.

Let us now discuss the evolution of the electronic band structure after K adsorption. A freshly K-doped surface with identical doping and spectroscopic features to the one corresponding to 0.09 ML of K was obtained after an annealing at 500 K of a K-saturated α -GeTe(111) surface. This point is discussed in Fig. 5 of the Appendix, in particular, by means of work function measurements. The electronic band structure of K/α -GeTe(111) along the $\bar{\Gamma}$ - \bar{K} high-symmetry direction [see Fig. 2(d)] exhibits sharp dispersions where bulk states are unchanged but the surface states SS_1 and SS_2 are rigidly shifted down to E_F with respect to the undoped surface. The upper part of the $SR_1 + B_1$ and $SR_2 + B_2$ branches is also slightly shifted to high binding energy as visible around $k = -0.1 \text{ \AA}^{-1}$. This is likely due to a small shift of the surface resonances which are now quasidegenerate with the bulk states close to E_F , in contrary to the α -GeTe(111) case [43]. In the case of bare α -GeTe(111), SS_1 and SS_2 , respectively, cross E_F at $k_{F,1} = -0.395 \text{ \AA}^{-1}$ and $k_{F,2} = -0.270 \text{ \AA}^{-1}$ whereas for K/α -GeTe(111), $k_{F,1} = -0.350 \text{ \AA}^{-1}$ and $k_{F,2} = -0.225 \text{ \AA}^{-1}$. This demonstrates that, for such a low coverage, the Rashba-type splitting is unaffected since the momentum splitting $\Delta k = k_{F,1} - k_{F,2}$ is not changed upon K doping. The energy shift of the Rashba split surface states towards higher binding energy is also visible from the Fermi

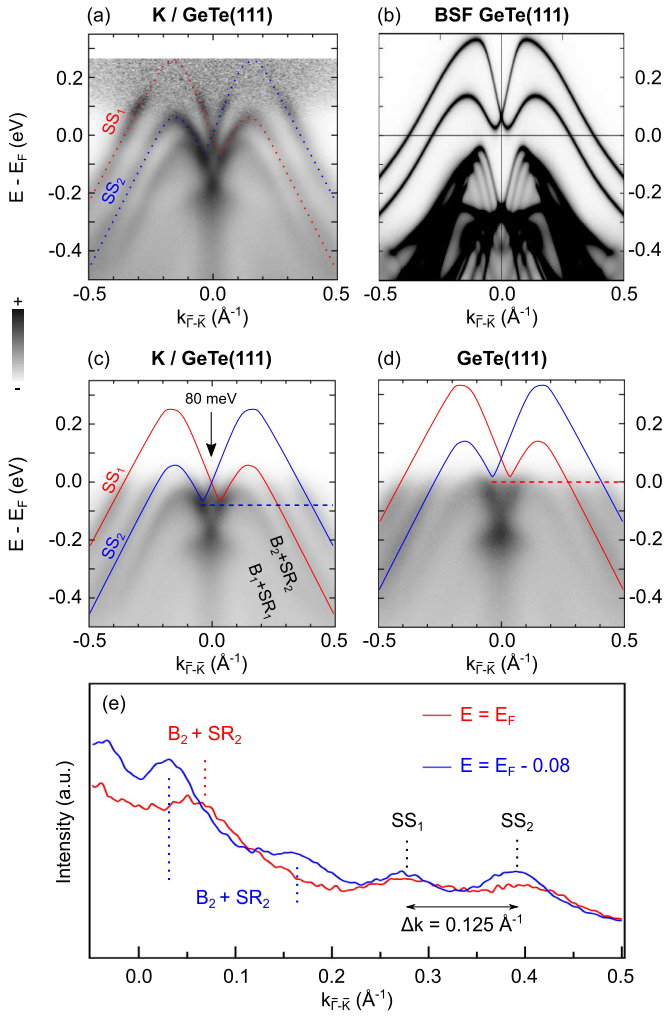


FIG. 3. (a) Fermi-Dirac distribution divided (at RT) ARPES spectrum along the $\bar{K}-\bar{\Gamma}-\bar{K}$ high-symmetry line of the 2D BZ for the $K/\alpha\text{-GeTe}(111)$ surface. The dashed-blue and red lines correspond to guides to the eye showing the full dispersions of SS_1 and SS_2 . (b) BSF calculations of Te-terminated $\alpha\text{-GeTe}(111)$. (c) and (d) ARPES spectra of $K/\alpha\text{-GeTe}(111)$ and $\alpha\text{-GeTe}(111)$. The solid red and blue lines are shifted by 80 meV up to E_F in (d) compared to panel (a). (e) Momentum distribution curves (MDCs) associated with the horizontal colored-dashed lines in panels (c) and (d).

surface [see Fig. 2(e)] where the two outer arcs are localized at lower momenta compared to the undoped surface, both along the $\bar{\Gamma}-\bar{K}$ and $\bar{\Gamma}-\bar{M}$ high-symmetry directions. Through K deposition, we, therefore, electron dope the $\alpha\text{-GeTe}(111)$ surface, leading to a rigid energy shift of the surface states down to the occupied states [44,45].

A direct consequence of the rigid energy shift of the surface electronic structure is the emergence of CP_{SS} at normal emission and E_F , related to pure surface states as highlighted by a black arrow in Fig. 2(d). It corresponds to the crossing of SS_1 and SS_2 which are degenerate at $k = 0$ as expected in a Rashba scenario due to the TRS conservation (Kramers degeneracy theorem). This CP_{SS} is not visible for the bare $\alpha\text{-GeTe}(111)$ because it occurs at too high energy in the unoccupied part of the band structure. We have plotted EDCs taken at the $\bar{\Gamma}$ point, i.e., at $k = 0$, of the bare and

K-doped $\alpha\text{-GeTe}(111)$ surfaces on the right-hand side of Figs. 2(a) and 2(d), respectively. In the case of $\alpha\text{-GeTe}(111)$, the EDC exhibits one dominant contribution, associated with the CP_{B+SR} of bulk and surface resonance states, also denoted by a black arrow in Fig. 2(a). In the K-doped case, there are two contributions: again the CP_{B+SR} but also the bottom of the surface states. The CP_{SS} is not distinguishable in the EDC because of the Fermi edge cutoff. As we will see below, it is possible to overcome this by dividing by the Fermi-Dirac distribution. Finally, the Fermi-surface map [see Fig. 2(e)] is accordingly modified: Some spectral weight appears exactly at the center of the snowflake at $\bar{\Gamma}$, forming a single point. Looking above E_F , e.g., at $E_F + 50$ meV [see Fig. 2(f)], this point becomes a ring corresponding to the lifting of degeneracy of the surface states away from the $\bar{\Gamma}$ point. Thus, the experimental band dispersions summarized in Fig. 2 indicate that, in contrast to earlier studies [34], surface states are not degenerate with bulk states at the $\bar{\Gamma}$ point. Moreover, the CP_{B+SR} are not affected by surface electronic doping because they remain located around 180 meV below E_F as previously reported [36].

Taking advantage of the well-defined band dispersions obtained at RT after K adsorption, we next divide the raw data of Fig. 2(d) by the Fermi-Dirac distribution, a well-known procedure used to gain access to the thermally occupied band structure in the range of a few tens of meV above E_F . It is particularly pertinent at RT because the width of the Fermi edge is approximately equal to 100 meV. Figure 3(a) displays the corresponding result. The detailed dispersion of the surface states is now evident, in particular, their CP at E_F and their linear dispersion in the vicinity of the $\bar{\Gamma}$ point. Guides to the eye are represented as red and blue dashed lines to highlight the two spin-polarized surface states SS_1 and SS_2 . In the present case, low-temperature measurements are not beneficial because they reduce our access to the thermally occupied part of the band structure (see Fig. 6 in the Appendix).

In order to obtain the energy position of the CP_{SS} in bare $\alpha\text{-GeTe}(111)$, Figs. 3(c) and 3(d) show ARPES spectra for $K/\alpha\text{-GeTe}(111)$ and the bare surface with guides to the eye, respectively, shifted by 0 and +80 meV in comparison to panel (a). For this size of rigid shift, the guides to the eye also fits very well to the surface states in the bare surface case. To confirm this, we have extracted MDCs for pertinent binding energies in Fig. 3(e). The MDCs are plotted at E_F for $\alpha\text{-GeTe}(111)$ (red line) and at 80 meV below E_F in the $K/\alpha\text{-GeTe}(111)$ case (blue line). They exhibit similar surface states contributions with a comparable Rashba-type splitting ($\Delta k = 0.125 \text{ \AA}^{-1}$) but different bulk/surface resonance ones. Indeed, whereas the MDC associated with the bare surface (red line) shows only one broad contribution centered at $k = +0.062 \text{ \AA}^{-1}$ for $B_2 + SR_2$, the one of the K-doped surface (blue line) exhibits two components, respectively, centered at $k = +0.030 \text{ \AA}^{-1}$ and $k = +0.165 \text{ \AA}^{-1}$. This not only demonstrates that upon K adsorption the bulk band structure is not modified, but also shows that the effect of K is to rigidly shift the surface states down to the E_F by the 80-meV value obtained by matching the momentum positions of SS_1 and SS_2 in the blue and red MDCs. Overall, our analysis allows us to obtain the energy position of the surface states in bare

α -GeTe(111) and conclude that their CP is positioned around 80 meV above E_F .

Finally, our observations are in excellent agreement with BSF calculations performed on a Te-terminated α -GeTe(111) surface as shown in Fig. 3(b). Indeed, the dispersion of the surface states is well reproduced by the simulations, especially their “V shape” above the CP at $\bar{\Gamma}$ and their nonparabolic character, remarkably different from nearly free electrons dispersion. The energy position of the CP_{SS} is theoretically estimated to be 70 meV above E_F , in very good agreement with our experimental findings.

IV. CONCLUSIONS

To summarize, direct measurement of the Rashba split surface states of α -GeTe(111) has been experimentally realized thanks to K doping, in excellent agreement with state-of-the-art band structure calculations. We report the detailed dispersion of these states in the vicinity of the $\bar{\Gamma}$ point, confirming the existence of their CP in the bulk band gap and, by consequence, the associated Rashba scenario with TRS conservation. We also present the definitive proof that the CP of the surface states is located in the unoccupied part of the band structure, a finding which was not properly established from previous experimental and theoretical works [34,37]. Furthermore, we show that the surface states shift in energy upon K doping but *not* the bulk states. By consequence, we can conclude that the shift of the surface states is not the result of band bending and that they are decoupled from the bulk states. We also find that K doping does not affect their Rashba splitting. These results clearly prove that the *intrinsic* origin of the giant Rashba splitting of the surface states of α -GeTe(111) is largely arising from the inversion symmetry breaking in the bulk. If it would be merely a surface effect, the Rashba splitting of the surface states should significantly change upon surface doping. Our paper is, therefore, an important step for the understanding of the electronic band structure of α -GeTe(111), which is a promising material for multifunctional spintronics devices. It opens the way for future investigations of the full spin texture of the surface states via spin-resolved ARPES around $\bar{\Gamma}$. Finally, our results call for additional studies of both the unoccupied band structure and the electron dynamics of the system *via* time-resolved ARPES measurements, in particular, the study of the potential relaxation channels of the electrons from the conduction band to the surface states.

ACKNOWLEDGMENTS

This project was supported by the Swiss National Science Foundation (SNSF) Grant No. P00P2_170597 and the Austrian Science Funds (FWF), Project No. P30860-N27. L.N. and J.M. would like to thank the CEDAMNF Project financed by the Ministry of Education, Youth and Sports of Czech Republic, Project No. CZ.02.1.01/0.0/0.0/15_003/0000358 and the Czech Science Foundation (GACR), Project No. 20-18725S. We are very grateful to P. Aebi for fruitful discussions and for sharing with us his photoemission setup. Skillful technical assistance was provided by F. Bourqui, B. Hediger, and O. Raetz.

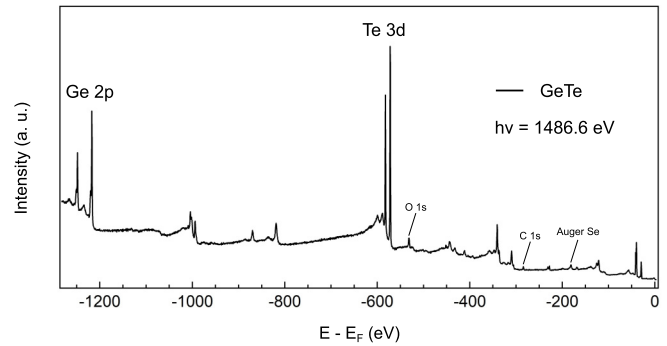


FIG. 4. XPS spectrum of the α -GeTe(111) surface obtained after the initial desorption procedure to remove the capping layer. The photon energy is $h\nu = 1486.6$ eV.

APPENDIX

1. X-ray photoemission spectroscopy

Figure 4 displays the XPS spectrum of the bare α -GeTe(111) surface, obtained after the desorption of the capping layer. It mainly exhibits the Ge 2*p* and Te 3*d* core levels with weak secondary electrons and contaminants signals, demonstrating the quality of the surface.

2. Extended data and sample characterization

Hereafter, we discuss how we have obtained the sample corresponding to the ARPES data on the K/α -GeTe(111) surface presented in the main text. We started from a bare α -GeTe(111) sample which has been characterized with LEED, ARPES, work function measurements, and XPS. The corresponding data are shown in Fig. 5. LEED, ARPES and XPS have been discussed in the main text. The work function measurement of the material is performed by applying a bias of -8 V and measuring the energy cutoff of the secondary electrons (E_{vac}). The work function of the material is given by the formula: $\Phi = h\nu - (E_F - E_{vac})$ with E_F defined as the kinetic energy of the electrons at the Fermi level. By performing such a measurement on α -GeTe(111), we extract a value of 4.7 eV, well above the K bulk value (1.8 eV). So, we are expecting an electron transfer from the K adatoms to the α -GeTe(111) surface.

After 1 min of potassium deposition (0.09 ML estimated using STM), the LEED pattern is more diffuse, and the ARPES data exhibit a shift of 80 meV of the surface states down to the Fermi level without any modification of the bulk states [Fig. 5(b)]. As expected, we observe a reduction of the work function down to 4.0 eV [Fig. 5(e)]. Increasing the amount of potassium at the surface by a factor of 10, i.e., 10 min of deposition, leads to the loss of both the LEED spots and the ARPES band structure [Fig. 5(c)], and to a decrease in the work function down to 3.4 eV.

By further annealing at 500 K, we recover the same LEED and ARPES data as those obtained after 1 min of deposition [Fig. 5(d)] and work function value of 4.0 eV. This demonstrates that the annealing procedure allows for recovering an electron doping level of the K-saturated α -GeTe(111) surface identical to the one obtained after 1 min of K deposition.

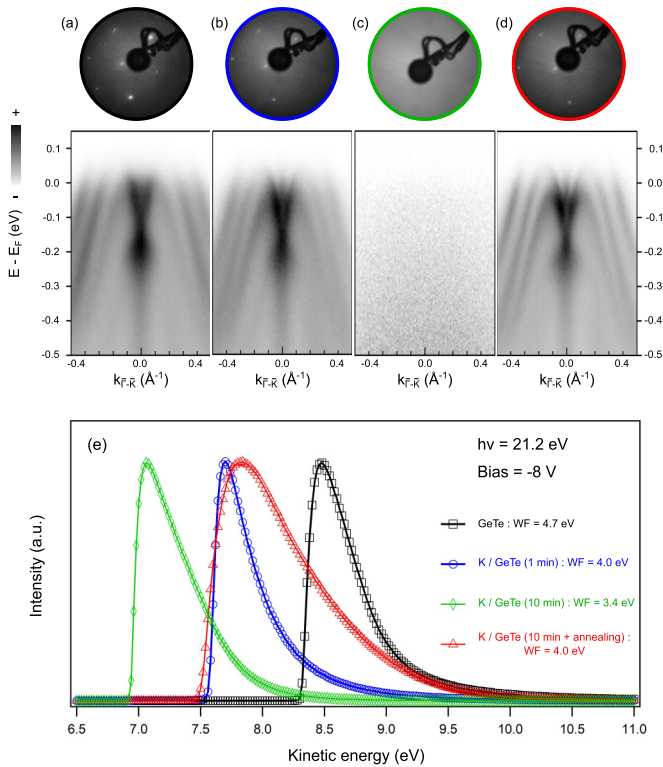


FIG. 5. Series of LEED patterns (top) and ARPES spectra (bottom) along the $\bar{K}-\bar{\Gamma}-\bar{K}$ high-symmetry line of the 2D BZ for (a) α -GeTe(111), (b) α -GeTe(111) with 1 min of K deposition, (c) α -GeTe(111) with 10 min of K deposition, and (d) α -GeTe(111) with 10 min of deposition followed by an annealing at 500 K. Note that the spectra from panels (a) and (d) have been measured with better statistics and with a smaller energy step compared to panels (b) and (c). (e) Corresponding low-energy cutoff of the secondary electrons as probed by HeI_α photoemission by applying a voltage bias of -8 V to the sample surface.

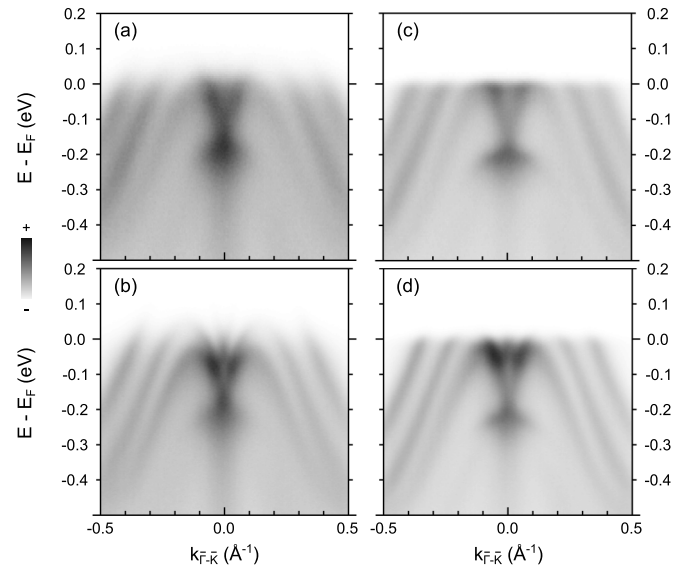


FIG. 6. Temperature-dependent ARPES measurements along the $\bar{K}-\bar{\Gamma}-\bar{K}$ high-symmetry line of the 2D BZ at RT for (a) α -GeTe(111) and (b) K/α -GeTe(111) and at $T = 80$ K for (c) α -GeTe(111) and (d) K/α -GeTe(111).

3. ARPES as a function of the temperature

To show the temperature dependence of the band structure of α -GeTe(111) and K/α -GeTe(111), Fig. 6 displays ARPES measurements taken at RT and at $T = 80$ K. As can be seen, the dispersions of the surface, bulk, and surface resonance states are the same. The only effects of lowering the temperature are to reduce both the spectral broadening and the width of the energy region in the thermally occupied part of the band structure that can be probed thanks to the thermal broadening of the Fermi edge.

- [1] Y. A. Bychkov and E. I. Rashba, *JETP Lett.* **39**, 78 (1984).
- [2] S. LaShell, B. A. McDougall, and E. Jensen, *Phys. Rev. Lett.* **77**, 3419 (1996).
- [3] M. Hoesch, M. Muntwiler, V. N. Petrov, M. Hengsberger, L. Patthey, M. Shi, M. Falub, T. Greber, and J. Osterwalder, *Phys. Rev. B* **69**, 241401(R) (2004).
- [4] E. Rotenberg, J. W. Chung, and S. D. Kevan, *Phys. Rev. Lett.* **82**, 4066 (1999).
- [5] M. Hochstrasser, J. G. Tobin, E. Rotenberg, and S. D. Kevan, *Phys. Rev. Lett.* **89**, 216802 (2002).
- [6] Y. M. Koroteev, G. Bihlmayer, J. E. Gayone, E. V. Chulkov, S. Blügel, P. M. Echenique, and P. Hofmann, *Phys. Rev. Lett.* **93**, 046403 (2004).
- [7] O. Krupin, G. Bihlmayer, K. Starke, S. Gorovikov, J. E. Prieto, K. Döbrich, S. Blügel, and G. Kaindl, *Phys. Rev. B* **71**, 201403(R) (2005).
- [8] K. Sugawara, T. Sato, S. Souma, T. Takahashi, M. Arai, and T. Sasaki, *Phys. Rev. Lett.* **96**, 046411 (2006).
- [9] H. Cercellier, C. Didiot, Y. Fagot-Revurat, B. Kierren, L. Moreau, D. Malterre, and F. Reinert, *Phys. Rev. B* **73**, 195413 (2006).
- [10] C. R. Ast, J. Henk, A. Ernst, L. Moreschini, M. C. Falub, D. Pacilé, P. Bruno, K. Kern, and M. Grioni, *Phys. Rev. Lett.* **98**, 186807 (2007).
- [11] T. Hirahara, K. Miyamoto, A. Kimura, Y. Niinuma, G. Bihlmayer, E. V. Chulkov, T. Nagao, I. Matsuda, S. Qiao, K. Shimada *et al.*, *New J. Phys.* **10**, 083038 (2008).
- [12] D. Hsieh, Y. Xia, L. Wray, D. Qian, A. Pal, J. H. Dil, J. Osterwalder, F. Meier, G. Bihlmayer, C. L. Kane *et al.*, *Science* **323**, 919 (2009).
- [13] J. H. Dil, *J. Phys.: Condens. Matter* **21**, 403001 (2009).
- [14] K. Sakamoto, H. Kakuta, K. Sugawara, K. Miyamoto, A. Kimura, T. Kuzumaki, N. Ueno, E. Annese, J. Fujii, A. Kodama *et al.*, *Phys. Rev. Lett.* **103**, 156801 (2009).
- [15] A. Varykhalov, D. Marchenko, M. R. Scholz, E. D. L. Rienks, T. K. Kim, G. Bihlmayer, J. Sánchez-Barriga, and O. Rader, *Phys. Rev. Lett.* **108**, 066804 (2012).
- [16] A. Tamai, W. Meevasana, P. D. C. King, C. W. Nicholson, A. de la Torre, E. Rozbicki, and F. Baumberger, *Phys. Rev. B* **87**, 075113 (2013).
- [17] E. E. Krasovskii, *J. Phys.: Condens. Matter* **27**, 493001 (2015).

- [18] P. Zhang, J.-Z. Ma, Y. Ishida, L.-X. Zhao, Q.-N. Xu, B.-Q. Lv, K. Yaji, G.-F. Chen, H.-M. Weng, X. Dai *et al.*, *Phys. Rev. Lett.* **118**, 046802 (2017).
- [19] K. Yaji, A. Harasawa, K. Kuroda, R. Li, B. Yan, F. Komori, and S. Shin, *Phys. Rev. B* **98**, 041404(R) (2018).
- [20] M. S. Bahramy, R. Arita, and N. Nagaosa, *Phys. Rev. B* **84**, 041202(R) (2011).
- [21] K. Ishizaka, M. S. Bahramy, H. Murakawa, M. Sakano, T. Shimojima, T. Sonobe, K. Koizumi, S. Shin, H. Miyahara, A. Kimura *et al.*, *Nat. Mater.* **10**, 521 (2011).
- [22] A. Crepaldi, L. Moreschini, G. Autes, C. Tournier-Colletta, S. Moser, N. Virk, H. Berger, P. Bugnon, Y. J. Chang, K. Kern *et al.*, *Phys. Rev. Lett.* **109**, 096803 (2012).
- [23] G. Landolt, S. V. Ereameev, Y. M. Koroteev, B. Slomski, S. Muff, T. Neupert, M. Kobayashi, V. N. Strocov, T. Schmitt, Z. S. Aliev *et al.*, *Phys. Rev. Lett.* **109**, 116403 (2012).
- [24] M. Sakano, J. Miyawaki, A. Chainani, Y. Takata, T. Sonobe, T. Shimojima, M. Oura, S. Shin, M. S. Bahramy, R. Arita *et al.*, *Phys. Rev. B* **86**, 085204 (2012).
- [25] G. Landolt, S. V. Ereameev, O. E. Tereshchenko, S. Muff, B. Slomski, K. A. Kokh, M. Kobayashi, T. Schmitt, V. N. Strocov, J. Osterwalder *et al.*, *New J. Phys.* **15**, 085022 (2013).
- [26] G. Landolt, S. V. Ereameev, O. E. Tereshchenko, S. Muff, K. A. Kokh, J. Osterwalder, E. V. Chulkov, and J. H. Dil, *Phys. Rev. B* **91**, 081201(R) (2015).
- [27] J. C. Rojas Sánchez, L. Vila, G. Desfonds, S. Gambarelli, J. P. Attané, J. M. De Teresa, C. Magén, and A. Fert, *Nat. Commun.* **4**, 2944 (2013).
- [28] E. Lesne, Y. Fu, S. Oyarzun, J. C. Rojas Sánchez, D. C. Vaz, H. Naganuma, G. Sicoli, J.-P. Attané, M. Jamet, E. Jacquet *et al.*, *Nat. Mater.* **15**, 1261 (2016).
- [29] S. Oyarzún, A. K. Nandy, F. Rortais, J. C. Rojas-Sánchez, M. T. Dau, P. Noël, P. Laczkowski, S. Pouget, H. Okuno, L. Vila *et al.*, *Nat. Commun.* **7**, 13857 (2016).
- [30] A. Soumyanarayanan, N. Reyren, A. Fert, and C. Panagopoulos, *Nature (London)* **539**, 509 (2016).
- [31] C. Rinaldi, J. C. Rojas Sánchez, R. N. Wang, Y. Fu, S. Oyarzun, L. Vila, S. Bertoli, M. Asa, L. Baldrati, M. Cantoni *et al.*, *APL Mater.* **4**, 032501 (2016).
- [32] D. Di Sante, P. Barone, R. Bertacco, and S. Picozzi, *Adv. Mater.* **25**, 509 (2013).
- [33] S. Picozzi, *Front. Phys.* **2**, 10 (2014).
- [34] M. Liebmann, C. Rinaldi, D. Di Sante, J. Kellner, C. Pauly, R. N. Wang, J. E. Boschker, A. Giussani, S. Bertoli, M. Cantoni, L. Baldrati, M. Asa, I. Vobornik, G. Panaccione, D. Marchenko, J. Sánchez-Bariga, O. Rader, R. Calarco, S. Picozzi, R. Bertacco, and M. Morgenstern, *Adv. Mater.* **28**, 560 (2016).
- [35] J. Krempaský, S. Muff, F. Bisti, M. Fanciulli, H. Volfová, A. P. Weber, N. Pilet, P. Warnicke, H. Ebert, J. Braun, F. Bertran, V. V. Volobuev, J. Minár, G. Springholz, J. H. Dil, and V. N. Strocov, *Nat. Commun.* **7**, 13071 (2016).
- [36] J. Krempaský, H. Volfová, S. Muff, N. Pilet, G. Landolt, M. Radović, M. Shi, D. Kriegner, V. Holý, J. Braun, H. Ebert, F. Bisti, V. A. Rogalev, V. N. Strocov, G. Springholz, J. Minár, and J. H. Dil, *Phys. Rev. B* **94**, 205111 (2016).
- [37] C. Rinaldi, S. Varotto, M. Asa, J. Ślawińska, J. Fujii, G. Vinai, S. Cecchi, D. Di Sante, R. Calarco, I. Vobornik, G. Panaccione, S. Picozzi, and R. Bertacco, *Nano Lett.* **18**, 2751 (2018).
- [38] J. Krempaský, S. Muff, J. Minar, N. Pilet, M. Fanciulli, A. P. Weber, E. B. Guedes, M. Caputo, E. Müller, V. V. Volobuev, M. Gmitra, C. A. F. Vaz, V. Scagnoli, G. Springholz, and J. H. Dil, *Phys. Rev. X* **8**, 021067 (2018).
- [39] G. Kremer, K. Zhu, T. Pierron, V. Fournée, J. Ledieu, S. Andrieu, B. Kierren, L. Moreau, D. Malterre, K. He, Q. K. Xue, Y. Fagot-Revurat, and Y. Lu, *J. Phys. D: Appl. Phys.* **52**, 494002 (2019).
- [40] H. Ebert, D. Ködderitzsch, and J. Minár, *Rep. Prog. Phys.* **74**, 096501 (2011).
- [41] P. Bauer Pereira, I. Sergueev, S. Gorsse, J. Dadda, E. Müller, and R. P. Hermann, *Phys. Status Solidi B* **250**, 1300 (2013).
- [42] V. L. Deringer, M. Lumeij, and R. Dronskowski, *J. Phys. Chem. C* **116**, 15801 (2012).
- [43] J. Krempaský, M. Fanciulli, L. Nicolaï, J. Minár, H. Volfová, O. Caha, V. V. Volobuev, J. Sánchez-Barriga, M. Gmitra, K. Yaji, K. Kuroda, S. Shin, F. Komori, G. Springholz, and J. H. Dil, *Phys. Rev. Research* **2**, 013107 (2020).
- [44] M. A. Hossain, J. D. F. Mottershead, D. Fournier, A. Bostwick, J. L. McChesney, E. Rotenberg, R. Liang, W. N. Hardy, G. A. Sawatzky, I. S. Elfimov *et al.*, *Nat. Phys.* **4**, 527 (2008).
- [45] T. J. Boyle, A. Rossi, M. Walker, P. Carlson, M. K. Miller, J. Zhao, P. Klavins, C. Jozwiak, A. Bostwick, E. Rotenberg *et al.*, *Phys. Rev. B* **100**, 081105(R) (2019).